Molecular structure of and exchange coupling in a bis(semiquinone) complex

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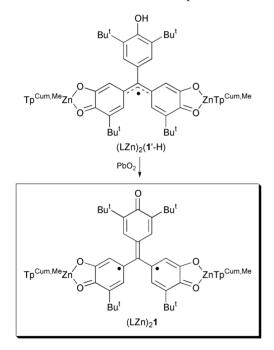
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The first example of a structurally characterized, rationally designed, triplet ground-state bis(semiquinone) complex is presented; the intramolecular exchange coupling is ferromagnetic with J = +209 cm⁻¹.

Molecular structures of organic biradical ligands are important to those concerned with construction of molecule-based magnetic materials using the metal-radical approach,^{1,2} and to those with an inherent interest in electronic properties of openshell organic molecules.³

Recently, we reported several semiquinone-type ligands to be used as building blocks for high-spin molecules and open-shell materials.^{4,5} Among these new ligands is S = 1/2 quinonemethide-semiquinone (QMSQ) present in the complex (LZn)₂(**1**'-H), the structure of which has been reported as well as its



conversion to both a mixed-valent species and to a biradical, $(LZn)_2\mathbf{1}$, shown below.^{5,6} Herein, we report both the molecular structure and magnetic properties of S = 1 (LZn)₂**1** [Tp^{Cum,Me} = hydrotris(3-cumenyl-5-methylpyrazolyl)borate].

For clarity, we will refer to the mono-oxygenated rings of $(LZn)_2(1'-H)$ and $(LZn)_21$ as A-rings, and the dioxygenated rings as either the dioxolene- or the B-rings. Structural changes that accompany conversion of $(LZn)_2(1'-H)$ to biradical $(LZn)_21$ are consistent with the Lewis structures above.

The crystal structure of biradical $(LZn)_2\mathbf{1}$ shows that the asymmetric unit contains two symmetry-independent molecules that differ primarily in torsion angles of A- and B-rings (and cumenyl ring torsions), and are therefore different conformers that also have an opposite sense of chirality.[†] An ORTEP of $(LZn)_2\mathbf{1}$ is shown in Fig. 1, where hydrogens and cumenyl

groups have been omitted for clarity. Bond lengths, angles, and torsions are summarized in Table 1 and Fig. 2.

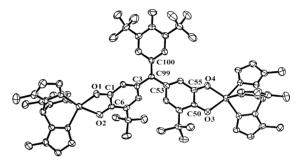


Fig. 1 50% ORTEP of (LZn)₂1. Hydrogen atoms, cumenyl groups and THF molecules have been omitted for clarity.

Bond	R _{av} /pm
C1-O1, C55-O4 C6-O2, C50-O3 C1-C6, C50-C55 C3-C99, C53-C99 C99-C100	129.35, 128.35 128.15, 127.15 146.05, 146.55 148.30, 146.95 137.65
Ring	Average torsion angles
C1–C6, C50–C55 C100–C105	50.33, 42.68 18.05

^{*a*} Averages for conformers I and II, the two symmetry-independent molecules in the asymmetric unit. ^{*b*} Torsion angles are relative to the plane made by atoms C3–C53–C99–C100.

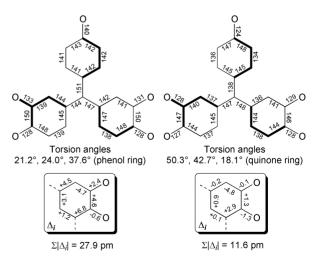


Fig. 2 Comparison of structural features of $(LZn)_2(1'-H)$ (left) and $(LZn)_21$ (right). Averaged bond lengths (in pm) are shown for $(LZn)_21$. See text for details.

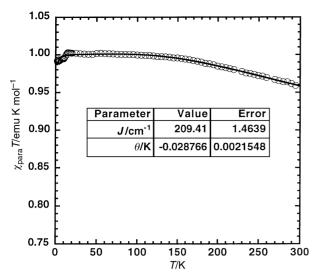


Fig. 3 Temperature dependence of $\chi_{para}T$ for biradical (LZn)₂1.

The A-ring C–O bond of $(LZn)_21$ (*ca.* 124 pm) is substantially shorter than the A-ring C–O single bond of $(LZn)_2(1'-H)$ (140 pm),⁵ consistent with carbonyl character for the former. In fact, as seen in Table 1 and Fig. 2, the pattern of long and short bond lengths for the A-ring in $(LZn)_21$ confirms its quinoidal character.

The A-ring of $(LZn)_2 \mathbf{1}$ is twisted 18° while the dioxolene rings are twisted *ca*. 46° from the plane made by atoms C3, C53, C99 and C100, allowing conjugation of the dioxolene rings with the A-ring.

Fig. 2 shows bond length deviations $(\Sigma |\Delta_i|)$ of the B-rings from a 3,5-di-*tert*-butylorthosemiquinone ring.^{7,8} As we pointed out previously, the dioxolene rings of S = 1/2 $(LZn)_2(1'-H)$ are quite different from those of a semiquinone $(\Sigma |\Delta_i| = 27.9 \text{ pm})$ due to the QMSQ bond formulation.⁵ However, $\Sigma |\Delta_i| = 11.6 \text{ pm}$ for the dioxolene rings of $(LZn)_2 1$ indicating a greater similarity to a semiquinone, but also consistent with some delocalization of spin/charge into the Aring, since we maintain that $\Sigma |\Delta_i| \neq 0$ is an indication of a delocalized electronic structure.

The magnetic susceptibility of microcrystalline $(LZn)_2 \mathbf{1}$ was measured from 2 to 300 K using a SQUID magnetometer, and is plotted as $\chi_{para}T$ vs. T in Fig. 3. Modeling the temperaturedependent $\chi_{para}T$ product of S = 1 molecules can be achieved by fitting to a modified Bleaney–Bowers expression:^{9,10}

$$\chi_{\text{para}}T = \frac{2Ng^2\beta^2T}{k(T-\vartheta\left[3 + \exp\left(\frac{-2J}{kT}\right)\right]}$$

where, g is the isotropic Landé constant (g = 2.0023), β is the Bohr magneton, T is temperature in Kelvin, k is Boltzmann's constant, ϑ is the Weiss constant representing a weak intermolecular interaction, and 2J is the singlet-triplet energy gap (J > 0 for triplet ground-state and J < 0 for singlet ground-state).

The best fit of the data give $J = +209 \pm 1 \text{ cm}^{-1}$ and $\vartheta \approx -0.03 \text{ K}$. The very small Weiss constant is consistent with negligible intermolecular interactions due to the large, insulating hydrotrispyrazolylborate ligand, while the moderately large, positive exchange parameter, J, indicates that $(LZn)_21$ is a triplet ground-state biradical consistent with its π -connectivity. The symmetry of the bis(semiquinone) ligand of $(LZn)_21$ precludes degeneracy of the SOMOs, nonetheless, delocalization of spin into the quinone-methide unit (A-ring) gives rise to a substantial exchange integral, stabilizing the triplet state relative to the singlet state. However, the different spin distribution, lower symmetry, and therefore attenuated delocalization into the A-ring results in a slightly smaller J for $(LZn)_21$ than for Yang's biradical $(J \ge +245 \text{ cm}^{-1})$.^{11,12}

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Notes and references

† *Crystal data* for (LZn)₂**1**: C₁₂₆H₁₆₀B₂N₁₂O_{8.25}Zn₂ which includes 3.25 molecules of THF solvate per Zn₂ complex, red crystals, M = 2127.02, triclinic, space group $P2_1/n$; a = 38.110(4), b = 18.827(2), c = 38.448(4)Å, $\beta = 118.546(2)^\circ$, V = 24233(4)Å³, $D_c = 1.166$ g cm⁻³, Z = 8, T = 158(2) K, R = 0.0730; wR2 = 0.2174, GOF = 1.016.

Measurements were carried out on a standard Siemens SMART CCDbased X-ray diffractometer equipped with a normal focus Mo-target X-ray tube ($\lambda = 0.71073$ Å) operated at 2000 W power (50 kV, 40 mA). The structure was solved and refined with the Brüker SHELXTL (version 5.10) software package (G. M. Sheldrick, Bruker AXS, Madison, WI, 1997). All non-hydrogen atoms were refined anisotropically with the hydrogen atoms placed in idealized positions. Modest restraints were placed on the lattice solvates to retain chemically reasonable geometries. Refinement was carried out using blocked least-squares matrices. F(000) = 9088; reflections collected, 184229; no. of unique reflections, 34986; no. of parameters, 2655; R = 0.0730; wR2 = 0.2174; ($\Delta \rho$)_{max} = 1.967 e Å⁻³.

CCDC 182/1860. See http://www.rsc.org/suppdata/cc/b0/b008037n/ for crystallographic files in .cif format.

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